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WISE/NEOWISE OBSERVATIONS OF COMET 103P/HARTLEY 2

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ABSTRACT

We report results based on mid-infrared photometry of comet 103P/Hartley 2 taken during 2010 May 4–13 (when the comet was at a heliocentric distance of 2.3 AU, and an observer distance of 2.0 AU) by the *Wide-field Infrared Survey Explorer*. Photometry of the coma at 22 μm and data from the University of Hawaii 2.2 m telescope obtained on 2010 May 22 provide constraints on the dust particle size distribution, $d \log n / d \log m$, yielding power-law slope values of $\alpha = -0.97 \pm 0.10$, steeper than that found for the inbound particle fluence during the Stardust encounter of comet 81P/Wild 2. The extracted nucleus signal at 12 μm is consistent with a body of average spherical radius of 0.6 ± 0.2 km (one standard deviation), assuming a beaming parameter of 1.2. The 4.6 μm band signal in excess of dust and nucleus reflected and thermal contributions may be attributed to carbon monoxide or carbon dioxide emission lines and provides limits and estimates of species production. Derived carbon dioxide coma production rates are $3.5(\pm 0.9) \times 10^{24}$ molecules per second. Analyses of the trail signal present in the stacked image with an effective exposure time of 158.4 s yields optical-depth values near 9×10^{-10} at a delta mean anomaly of 0.2 deg trailing the comet nucleus, in both 12 and 22 μm bands. A minimum chi-squared analysis of the dust trail position yields a beta-parameter value of 1.0×10^{-4} , consistent with a derived mean trail-grain diameter of 1.1/ ρ cm for grains of ρ g cm $^{-3}$ density. This leads to a total detected trail mass of at least 4×10^{10} ρ kg.

Key words: comets: individual (103P/Hartley 2) – infrared: planetary systems

1. INTRODUCTION

In our solar system, only five cometary bodies have been studied up close via spacecraft encounters. A careful characterization of these few bodies and their behavior over time therefore provides vital ground-truth tests to the interpretations drawn for the larger data set of cometary bodies obtained from telescopic observations. The *EPOXI/DIXI* encounter of 103P/Hartley 2 (103P) provided a unique view of a Jupiter-family comet (JFC) near its perihelion, when it was most active. The extended ground-based monitoring of 103P (Meech et al. 2011) indicated that activity increased over the course of its much-studied perihelion approach. Here, we present a unique collection of combined data during its more sparsely sampled pre-perihelion period in the spring of 2010, over 5 months prior to the *EPOXI* encounter, when the comet was at nearly twice the distance from the Sun and was imaged by the *Wide-field Infrared Survey Explorer* mission (WISE; Wright et al. 2010).

The WISE mission surveyed the sky at four mid-IR wavelengths simultaneously, 3.35 μm (W1), 4.60 μm (W2), 11.56 μm (W3), and 22.08 μm (W4), with approximately one hundred times improved sensitivity over the *IRAS* mission (Wright et al. 2010). As part of an enhancement to the WISE data processing system called “NEOWISE,” funded by NASA’s Planetary Science Division, the WISE Moving Object Process-

ing Software (WMOPS) was developed to find solar system bodies in the WISE images. WMOPS successfully found a wide array of primitive bodies, including Near-Earth Objects (NEOs), Main Belt asteroids, comets, Trojans, and Centaurs. As of 2011 January, NEOWISE has identified more than 157,000 small bodies, including 123 comets (Mainzer et al. 2011a). Such infrared observations are useful for determining size and albedo distributions, thermo-physical related properties such as thermal inertia, the magnitude of non-gravitational forces, and surface roughness (Mainzer et al. 2011b). The subset of these bodies exhibiting cometary activity requires special treatment in the interpretation of such observations, owing to the material surrounding the solid nucleus, i.e., the contribution to the IR flux from the gas and dust, and to the variable nature of the observed brightness of the object attributable to outbursts. These IR imaging data provide unique opportunities to characterize four main components of comets: the coma dust and gas, the nucleus, and the extended dust trail. The subject of this paper, 103P, was of particular interest owing to the *EPOXI* spacecraft flyby scheduled in 2010 November. As a result, much supporting data, including visual-band wavelength observations near the time of the WISE observations, and a well-characterized light curve at similar distances, enhance the value of the WISE data. The mid-IR data in turn augment the more comprehensive picture of the *EPOXI/DIXI* mission target by providing a detailed view of its

Table 1
Mid-IR and Ground-based Observations of 103P

Date/Time (UT)	Telescope/Instrument ^a	Wavelength (μm)	Images	Exposures ^b	Comments
2010 Apr 17, 12:02	Pal. 200", LFC	0.7	3	60	Poor seeing, thin cirrus
2010 Apr 18, 10:03	SOAR 4.1 m, SOI	0.7	3	60	Photometric
2010 May 23, 13:52	UH 2.2 m, Tek2048	0.7	8	300	Photometric
2010 Jul 13, 09:30	Pal. 200", LFC	0.7	4	30	Photometric
2010 Aug 17, 03:56	SOAR 4.1 m, SOI	0.7	5	180	Photometric
2010 Sep 03, 06:46	Pal. 200", LFC	0.7	3	60	Photometric
2010 Nov 07, 09:35	Pal. 200", LFC	0.7	22	20	Photometric
2011 Jan 27, 05:58	Steward 2.3 m, 90 prime	0.7	1	15	Thin cirrus
2010 May 10, 08:18–2010 May 11, 11:59 ^c	<i>WISE</i>	3.4, 4.6, 11.6, 22.1	18	8.8	Coma and nucleus
2010 May 04, 06:13–2010 May 13, 04:33	<i>WISE</i>	11.6, 22.1	113	8.8	Debris Trail observations

Notes.

^a Pal 200": the Palomar 200" telescope, Mt. Palomar, CA, Southern Astronomical Research Telescope, Cerro Pachon, Chile, UH 2.2 m: the University of Hawaii 2.2 m telescope, Steward 2.3 m: the Steward Observatory 2.3 m telescope on Kitt Peak, with the 90 prime One camera.

^b Average exposure time in seconds.

^c Individual *WISE* exposure modified Julian Date values: 55326.34566, 55326.47797, 55326.61027, 55326.74258, 55326.87488, 55326.94097, 55327.00718, 55327.07327, 55327.13949, 55327.20558, 55327.27179, 55327.33788, 55327.40410, 55327.47018, 55327.60249, 55327.73479, 55327.86710, 55327.99940.

behavior at an earlier epoch. The *WISE* observations occurred at a heliocentric distance over twice that at the time of the 2010 November flyby. 103P also proved an excellent test case for the array of information the *WISE* data set has provided for the total set of comets.

2. OBSERVATIONS

In anticipation of the observation of multiple targets of interest, we applied for telescope time at several sites to obtain visual wavelength imaging and photometry near the times of the *WISE* observations. A combined campaign of ground-based and space-based telescopes has provided characterization of 103P from the previous perihelion passage through the time of the encounter and beyond (cf. Meech et al. 2011). Our team concentrated our telescope allocation time requests around the expected dates for the observation of 103P, but continued to observe 103P beyond that period as well. Table 1 lists the times of observations at the various ground-based sites, in addition to the *WISE* observations obtained on 2011 May 10–11 (see Figure 1). A total of eight nights of observation are listed, characterizing the behavior of the comet both at the time of the *WISE* observations and near the comet's perihelion and encounter times (2010 October 20 and November 4, UT, respectively).

A total of 18 four-color *WISE* images (Figure 1) were obtained of the nucleus and coma over the course of 39.7 hr. That this was more than the average 10 *WISE* images obtained per sky object was attributable to the largely pro-grade motion of the comet ($45\text{--}70\text{ arcsec hr}^{-1}$) moving in the direction of the scan progression of the spacecraft (*WISE* covers all ecliptic latitudes each day in a narrow band of sky at $90 \pm 2\text{ deg}$ elongation and uses the spacecraft orbital motion around the Sun to scan this band across all ecliptic latitudes over one year). Variation in the sky-motion rates were attributable mostly to the spacecraft's orbital motion; in any case this reflex motion created a maximum blurring of $\sim 0.02\text{ arcsec}$, an insignificant factor in the imaging, as the blur was $<1\%$ of the point-spread function (PSF) full width at half-maximum (FWHM) in the shortest wavelength ($2.75\text{ arcsec pixel}^{-1}$ average scale of the images and $5.5\text{ arcsec pixel}^{-1}$ in the 2×2 binned $22\text{ }\mu\text{m}$ images; Wright et al. 2010). The long, extended debris trail was imaged

piecewise over the course of nine days, using a total of 113 *WISE* images.

3. ANALYSIS

3.1. Optical Photometry

For all but the night of 2010 April 17 and 2011 January 27 the Palomar, SOAR, Mauna Kea, and Steward observatory ground-based observing conditions for 103P were photometric, and provided magnitudes to within a few percent uncertainty. After acquisition, the ground-based images were bias-subtracted and flat-field corrected. The *R*-band photometry obtained from the ground-based observations is summarized in Table 2, including flux magnitudes and geometry corrected $Af\rho$ dust column values (A'Hearn et al. 1984) as well as magnitudes. The fixed 2.5 and 5 arcsec aperture radius values listed in Table 2 are chosen to highlight the activity near the comet's coma center. The values represent the average of the aperture signals in the individual images, but special analysis was performed for the 2010 April 18 and May 23 (Figure 2) observations. Larger aperture radius values of 11 and 22 arcsec were used to match the aperture sizes necessary to obtain the signal from the poorest resolution *WISE* bands. We stacked and median combined the May 23rd UH frames into a single image to reduce the contribution of background stars to the large aperture and background annulus signals. Photometric analysis yielded a reflected-light signal in these large aperture values for comparison to the *WISE* observations by adjusting the 2010 April 18 and May 23 aperture magnitudes to the May 10th *WISE* spacecraft-to-comet and comet heliocentric distance. Interpolation between the adjusted 2010 April 18 and May 23 magnitudes yielded 2010 May 10 magnitudes of 18.45 ± 0.02 , and 18.11 ± 0.02 , and $Af\rho$ values of 0.85 and 0.69, for the 11 and 22 arcsec aperture signal, respectively. Note that the interpolation correction was 0.04 mag, within the photometric uncertainty of the April 18 magnitudes and on the order of what might be expected for the change in viewing phase from 26.1 to 27.3 deg for a cometary body (with slope parameters between 0.02 and $0.035\text{ mag deg}^{-1}$; Meech & Jewitt 1987). A single *V*-band 300 s exposure was taken the night of the 2010 May 23 observations, yielding *V* – *R* colors of 0.38 ± 0.03 , or nearly solar.

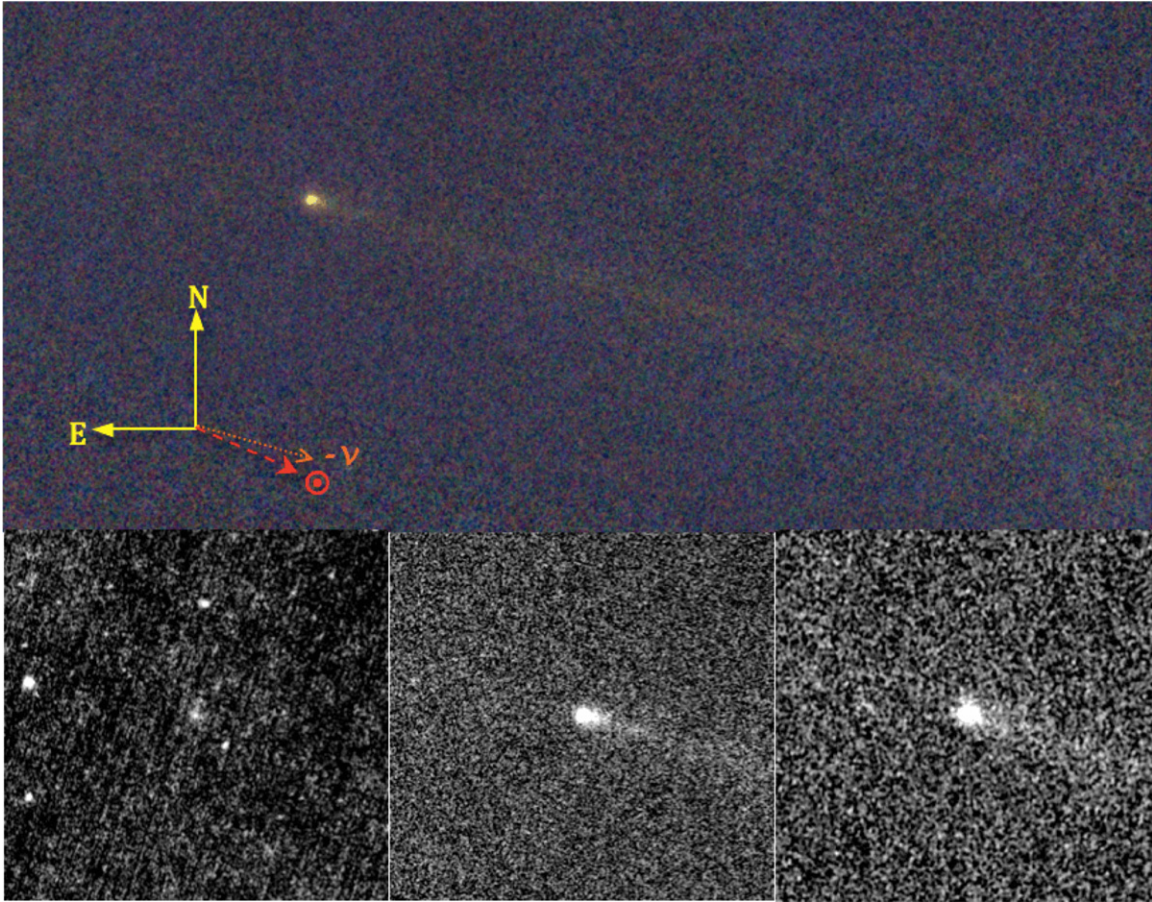


Figure 1. Three-color composite image of 103P/Hartley 2 from the *WISE* data taken on 2010 May 10–11. The 18 stacked 4.6, 11.6, and 22.1 μm images were mapped to blue, red, and green channels 22 arcmin on a side, insets of the individual band images close to the comet, 5 arcmin on a side, are shown below (W2, W3, and W4 images, left to right). The sky-projected anti-solar vector is indicated by the red dashed arrow, and the projected anti-velocity vector by the orange dotted arrow.

Table 2
R-band Optical Photometry

Date/Time (UT)	Airmass	$m_R, r_{\text{ap}} = 2''.5$	$m_R, r_{\text{ap}} = 5''$	$Af\rho, 2''.5$ (log cm)	$Af\rho, 5''$ (log cm)	R_{helio}/Δ
Apr 17, 12:02	1.70	19.90 ± 0.12	19.49 ± 0.19	1.08	0.94	2.49/2.52
Apr 18, 10:03	1.20	20.03 ± 0.04	19.62 ± 0.06	1.02	0.89	2.47/2.49
May 23, 13:52	1.12	19.46 ± 0.01	19.15 ± 0.01	1.00	0.82	2.19/1.81
Jul 13, 09:30	1.08	17.45 ± 0.02	16.87 ± 0.02	1.32	0.98	1.73/0.94
Aug 17, 03:56	2.46	15.99 ± 0.01	15.27 ± 0.01	1.49	1.05	1.43/0.52
Sep 3, 06:46	1.03	15.26 ± 0.02	14.54 ± 0.01	1.54	0.99	1.29/0.38
Nov 7, 09:35	1.44	13.60 ± 0.03	12.55 ± 0.01	1.70	1.02	1.07/0.17
Jan 27, 05:58	1.34	15.28 ± 0.20	14.63 ± 0.20	1.98	1.94	1.59/0.69

3.2. WISE Observations

The coma images were identified using the method described in the *WISE* archive explanatory supplement (<http://wise2.ipac.caltech.edu/docs/release/prelim/expsup/>), utilizing the WMOPS detections in combination with the single-frame (level 1b) archive source search. An additional frame with a near-edge detection (Modified Julian Date 55327.47018; see Table 1) was found from individually searching the archive scans spanning the dates of coverage (2010 May 10–11).

The *WISE* data were processed through the *WISE* Science Data System (WSDS) pipeline (Wright et al. 2010), to remove detector artifacts as well as bias and flat correct each image. The resultant images were stacked into the comet-centered images shown in Figure 1, using the WSDS co-adder routine specially

adapted to stack according to the comet’s rate of apparent sky motion. Aperture photometry was performed on the stacked images of each band for apertures using radii of 6, 9, 11, and 22 arcsec. Counts were converted to fluxes using the band-appropriate magnitude zero-points and zeroth-magnitude flux values provided in Wright et al. (2010), and an iterative fitting to a blackbody curve was conducted on the two long-wavelength bands to determine the appropriate color correction as listed in the same. Figure 3 shows the blackbody fit and the signal in W2, W3, and W4. The W1 image did not yield significant signal in any of the apertures. However, the reflected light signal was constrained by the visual wavelength data also shown on the graph. In Figure 3 fluxes, the flux contribution from the nucleus (see Figure 4) has been removed. Even so, note that the flux in W2 is significantly greater than the total dust reflected light and

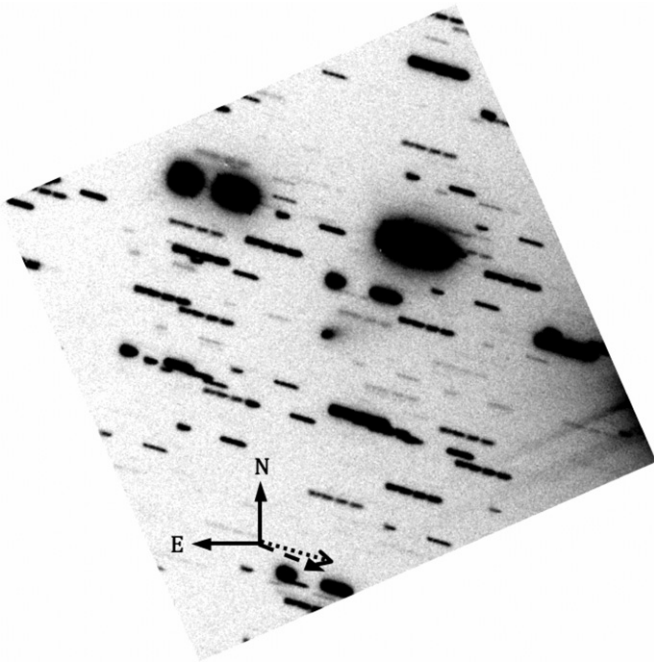


Figure 2. Optical image taken at the UH 2.2 m on the night of 2010 May 23. The image was taken with the Tek 2048 camera, $22'' \text{ pixel}^{-1}$, and shown here with 3 arcmin on a side. The sky-projected anti-solar vector is indicated by the dashed arrow, and the projected anti-velocity vector by the dotted arrow.

thermal contribution, suggesting another narrowband source of emission in the $4\text{--}5 \mu\text{m}$ range of the W2 passband.

In order to extract the nucleus signal, we used routines developed by our team (Lisse et al. 1999 and Fernández et al. 2000) to fit the coma as a function of angular distance from the central brightness peak along separate azimuths. The extracted nucleus signals in W3 and W4, $0.7 \pm 0.2 \text{ mJy}$ and $2.6 \pm 0.7 \text{ mJy}$, respectively, were fit to a NEATM model (Harris 1989, Delbo et al. 2003, and Mainzer et al. 2011b) with free and fixed beaming (η) parameters (Figure 4). Fits with η values fixed to 2.0, 1.2 (Stansberry et al. 2007), and 0.94 (Fernández et al. 2008) yielded diameters of 1.5, 1.1, and $0.9 (\pm 0.2) \text{ km}$. The geometric mean radius, based on radar measurements (Harmon et al. 2010), is closer to 0.58 km , or a diameter of 1.16 , more consistent with the $\eta = 1.2$ fixed value, and are closer to the 1.14 km diameter value derived from *Spitzer* measurements of the 2008 August 13, $r_h = 5.4 \text{ AU}$, $22 \mu\text{m}$ flux obtained by Lisse et al. (2009). The $\eta = 1.2$ fixed fit yields an albedo of $p_v = 0.038 \pm 0.004$, assuming an absolute magnitude of 18.8 consistent with a 0.6 km mean radius (Lisse et al. 2009). Note that each fit iteration requires an interpolation for surface temperature in the *WISE* bands (Wright et al. 2010), so that different flux values are derived for each final fit, as shown in Figure 4. The interpolated corrections for temperature are largest in W3. For decreasing η , the temperature increases. Note that the flux contribution from the nucleus in W2 was calculated to be less than 10^{-5} Jy , or significantly less than the uncertainty in the $4.6 \mu\text{m}$ flux. The W3 and W4 flux contributions from the nucleus were also subtracted from the total fluxes used to calculate the particle size distribution (PSD) shown in Figure 5.

3.3. Dust Trail

Observations of the 103P/Hartley 2 debris trail spanned the time period from May 4:13:35 to May 13:33:47, and a position on the sky of $1:3$ leading to 4° trailing the comet, covering a range from $-0:32$ to $+1:08$ in delta Mean Anomaly (dMA).

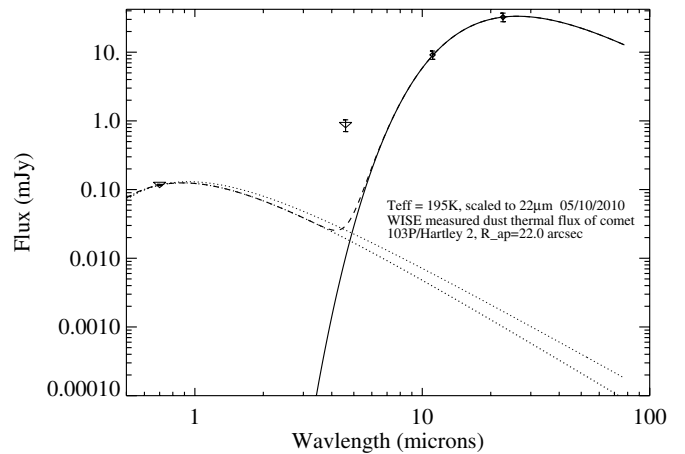


Figure 3. Coma temperature fit to the 22 arcsec aperture thermal photometry in the two longest *WISE* wavelength bands. Nucleus and reflected light signal were subtracted from the flux in the two bands prior to fitting. The $4.6 \mu\text{m}$ band coma signal is also shown, along with the *R*-band brightness measured on May 23rd, adjusted for distance, and with a small 4% adjustment for the projected activity based on multiple observations (see the text). Two reflected-light models are shown, one with a neutral reflectance (heavy dotted) and one with a reddening law based on Jewitt & Meech (1986) averaged out to $3.5 \mu\text{m}$ (light dotted line), i.e., with flux proportional to the wavelength to the approximately 0.2 power. The uncertainties to the temperature fit are on the order of $\pm 9 \text{ K}$, and the fitted temperature (195 K) closely matches the black body temperature for that distance (185 K). Excess W2 flux cannot be explained by thermal and reflected-light contributions of dust.

The debris trail of 103P is a low contrast feature with respect to the infrared background and requires co-addition to obtain reasonable signal-to-noise. The co-addition should be of the same trail features, for example, at the same dMA, to preserve any variations in the emission history.

W3 and W4 images were selected from the set of 113 *WISE* level 1b frames containing the trail, and the W3 images rebinned to the W4 pixel size. The background radiance of each image was fit with a two-dimensional low-order polynomial and subtracted from the image to reduce the effects of global radiance slopes. The resulting images were then shifted to bring the desired dMA to the center of a stacked (co-added) image of $508 \text{ columns} \times 200 \text{ rows}$, and rotated about that dMA to bring the projected orbit of 103P parallel to the stacked image rows. The co-added image thus contains the orbit along row 100 with the selected dMA at the center. The increment of dMA used was $0:04$, and the number of images stacked at each dMA varied from 14 to 19. Figure 6(a) shows co-added images for a stack of W3 and W4 frames centered at $\text{dMA} = 0:2$.

Radiance profiles normal to the trail were constructed from the mean and standard deviation of each row. Outlier rejection (of stars, instrumental artifacts, cosmic rays, etc.) was applied to each row using an iterative technique based on the theory of order statistics (Mandell 1964). The radiance profiles were fit to a Moffat (1969) function with an exponent of 0.5 to find the peak radiance, the position of the peak, and the FWHM. Radiance profiles derived from the two images shown in Figure 6(a) are plotted in Figure 6(b). The Moffat function, which is a modified Lorentz function, was chosen since it fits the peak, position, and shape of the radiance profile more accurately than the conventional Gaussian. The Lorentz form of the trail profile was originally suggested as a result of dynamical modeling of trail grains (D. Lien 1990, private communication). Data numbers (DNs) were converted to fluxes using the band-appropriate magnitude zero-points and zero-magnitude flux values provided

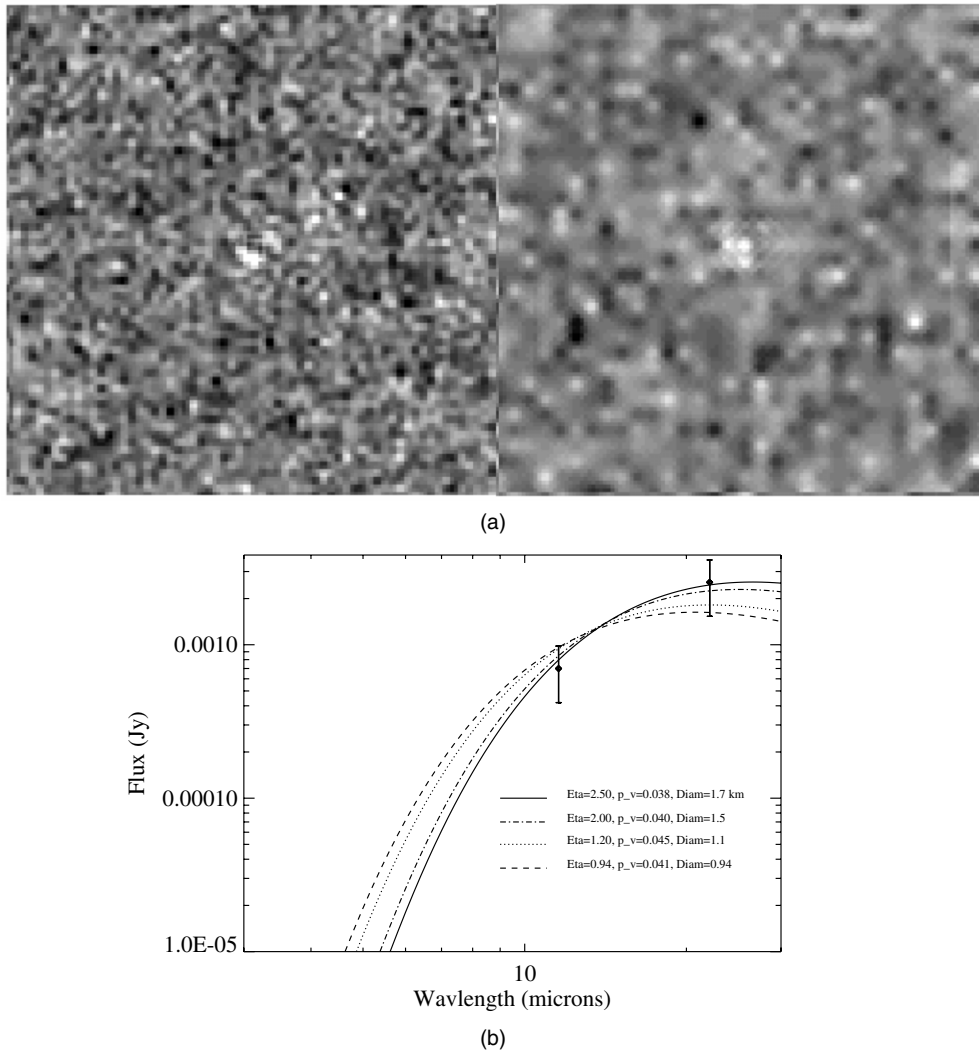


Figure 4. Thermal models of extracted nucleus signal. Figure 4(a) shows the W3 (left) and W4 (right) residual signal after coma extraction, using $1/r$ model coma fits. Figure 4(b) shows the thermal fits to the data. Unconstrained fits for beaming parameter (η) values yield fits of $\eta \sim 2.5$ (solid line), and diameter values of 1.7 km. η parameter values constrained to 2.0 (dot-dashed line), 1.2 (dashed line), and 0.94 (dotted line) yield diameters of 1.5, 1.1, and 0.9 km, respectively.

in Wright et al. (2010). The peak profile in-band radiances are plotted as a function of dMA in Figure 7. In this and subsequent figures (Figures 8 and 9), vertical dotted lines delineate a region in which the trail radiances are influenced by the coma and tail radiances. This is due, in part, to the small angle between the sky-projected anti-solar vector and the projected anti-velocity vector (see Figure 1). However, it is mainly a result of the “boxcar-like” averaging effect of taking row means. It is also responsible for the peak radiance occurring at dMA = 0.16 rather than at the nucleus (dMA = 0) since the trail is brighter behind than in front of the comet. The mean radiance of the trail leading the comet is 0.05 MJy sr⁻¹ in W4 and 0.02 MJy sr⁻¹ in W3. Behind the comet, the trail peak radiance falls from 0.14 MJy sr⁻¹ at 0°:2 dMA to 0.06 MJy sr⁻¹ at 1°:04 dMA in W4 and 0.06 MJy sr⁻¹ to 0.04 MJy sr⁻¹ in W3.

4. DISCUSSION

The three-color composite image and individual band images in Figure 1 show many of the summary features of the data set that provide unique characterization of 103P. The composite image shows an extended trail with thermal emission long-ward of 10 μ m and with little or no reflected-light component in the shorter wavelength 3.4 and 4.6 μ m bands. The trail also seems

to be noticeably missing from the deep-exposure, stacked UH 2.2 m image in Figure 2 taken within 12 days of the *WISE* data. There is also the hint of a dust trail in the 11.6 and 22.1 μ m bands that precedes the comet, indicative of large grains remaining from the comet’s previous passage. Note in the images that Hartley 2’s trail-ward brightness variations do not persist in the stacked subsets of the images and are likely caused by noise fluctuations rather than large-scale ejecta. At 4.6 μ m, the image is considerably compact, even when compared with the *R*-band image from the UH 2.2 m, suggesting that the dominant source of the brightness is not the same source as the reflected light or thermal emission of the dust, as is likely in the *WISE* bands W3 and W4, and the *R* band.

The aperture photometry at 11.6 and 22.1 μ m, in combination with the *R*-band photometry, provides constraints on the dust PSD. Proceeding as in Bauer et al. (2008), if we assume, starting with the longest wavelengths, that the thermal emission comes primarily from the dust particles near the size of the wavelength of emission, we can obtain an estimate of the number of large particles that fall within the aperture by scaling to the nucleus signal using the formula

$$n_g = (R/a)^2 \times (F_{\text{total},\lambda} - F_{\text{nuc},\lambda})/F_{\text{nuc},\lambda}, \quad (1)$$

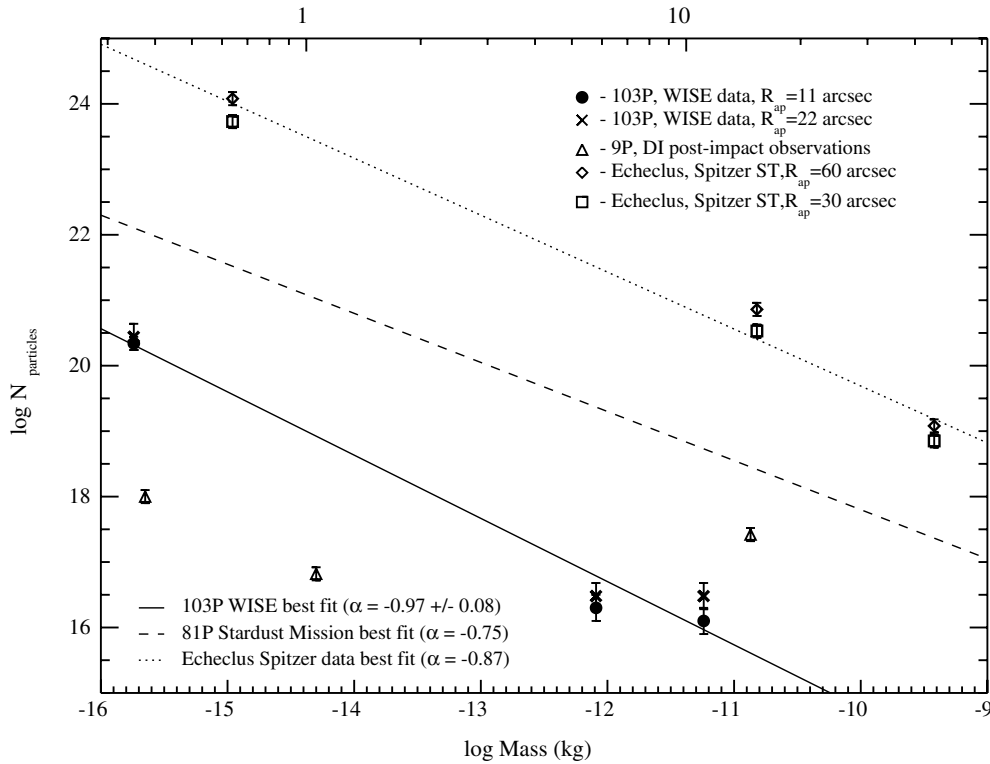


Figure 5. Particle size distribution (PSD). Log number is shown vertically, while log mass is shown on the bottom scale and the corresponding grain radius size, in μm , is shown on the scale above. The 103P data derived number of particles in the 11 arcsec aperture radius (pentagon) and 22 arcsec aperture (cross) are shown. The flux contributions from the nucleus have been subtracted. For comparison, *Deep Impact* particle densities (triangles), and Echeclus particle numbers (diamonds and squares, as presented in Bauer et al. 2008) are also shown. Stardust PSD slope ($\alpha = -0.75$, in $\log N/\log \text{kg}$ units, where N is the estimated total number of dust grains in the aperture) is shown as the dashed line, rescaled from dust fluence values to an aperture encompassing a similar ρ size. Echeclus' PSD best fit ($\alpha = -0.87$) is shown as a dotted line, and the solid line is the best fit to 103P PSD data ($\alpha = -0.97$).

Table 3
Coma Mid-IR Fluxes (mJy)

R_{Aperture}	$4.6 \mu\text{m}$	$12 \mu\text{m}$	$22 \mu\text{m}$
6''	0.16 ± 0.02	2.2 ± 0.3	5 ± 1
9''	0.29 ± 0.03	4.2 ± 0.6	11 ± 2
11''	0.39 ± 0.04	5.4 ± 0.7	15 ± 2
22''	0.86 ± 0.05	10 ± 1	31 ± 5

where n_g is the total count of grains of radius a (usually equal to $\frac{1}{2}$ the observing wavelength λ) and $F_{\text{total},\lambda}$ is the total flux, including the nucleus, while $F_{\text{nuc},\lambda}$ is the flux contribution from the nucleus. For the shorter thermal wavelengths, the same calculation was repeated, but the contribution to the flux from the larger particles, as estimated from the derived number of particles and projected area for the assumed particle sizes, is subtracted off. The reflected light contribution from the larger particles is also subtracted from the short-wavelength flux before deriving a similar particle number from

$$n_g = ((Af\rho \times \rho/p)/a)^2, \quad (2)$$

where $Af\rho$ and ρ are as defined in A'Hearn et al. (1984), p is the estimated particle reflectance (here 0.03), and a is as in Equation (1). The number of particles as a function of particle mass is shown in Figure 5 log-log plot, while derived values are shown in Table 3. Production rates for the dust sizes were calculated from the crossing time of the projected distance of the aperture (i.e., ρ , which is 15,994 and 31,889 km for 11 and 22 arcsec, respectively), assuming an ejection velocity $\sim 1 \text{ km s}^{-1}$.

Note that the dust production rate,

$$Q_{\text{dust}} = n_g \times m_g \times \rho/v_g, \quad (3)$$

is inversely proportional to the grain velocity, v_g , which is estimated here as similar to the CO_2 gas ejection velocity at this distance (cf. Pittichova et al. 2008). The number of coma dust grains, the dust mass estimates, and production rates as a function of size are summarized in Table 4. $\log m_g$ versus $\log n_g$ is shown in Figure 5 and provides approximate PSD power-law slopes (α) for the dust particle size distribution (dust PSD). For comparison, dust PSD data are shown from various measured cometary bodies, including the *Deep Impact* ejecta from the impactor experiment on 9P/Tempel 1, the active Centaur Echeclus, and the Stardust Mission target 81P/Wild 2. Our estimate of dust PSD for 103P yields a moderately steeper slope ($\alpha = -0.97 \pm 0.08$) when compared to these other bodies, suggesting, when compared to the smaller grains, relatively fewer large-grained dust particles were generated while the comet was at a distance of 2.3 AU. Furthermore, the 11.6 and 22.1 μm data points suggest that the 5.7 μm radius-scale grains are fewer than predicted by the simple power law. This may be indicative of a bifurcation in the distribution, where the mass is more strongly weighted to grain sizes greater than radii 5.7 μm , and may be consistent with the strong trail signal seen for 103P. Note that alternatively in the literature α is expressed as a size-dependent slope, rather than mass-dependent (cf. Fulle et al. 2004). The value quoted above translates into an $\alpha_{\text{size}} \approx 3.0 \pm 0.3$ ($\log N/\log a$), not an uncommon value for JFC comets that exhibit strong dust trails and weak silicate features, both indicative of abundant large dust grains (Lisse et al. 1998; Fulle

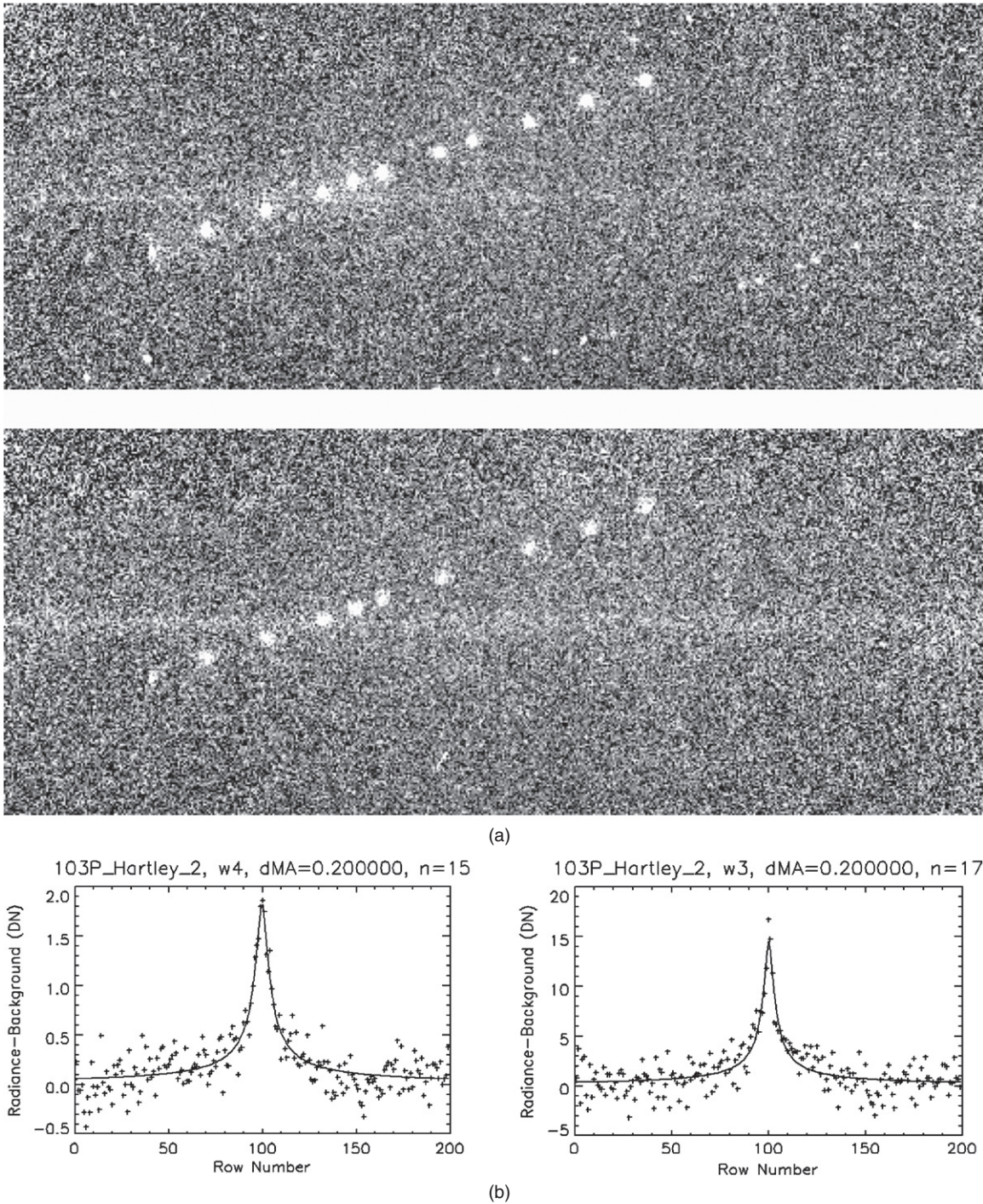


Figure 6. (a) W3 (top) and W4 (bottom) images shifted, rotated, and stacked with respect to delta mean anomaly and centered on $dMA = 0.2$. The W3 image is a stack of 17 frames, while the W4 image is a stack of 15 frames. The trail is seen in both images as a faint diffuse horizontal line through the center of the images. Note that the background stars running diagonally through the trail image. (b) Radiance profiles for the stacked W3 (left) and W4 (right) images shown in panel (a). Plotted points are the mean radiances for each row. The units are data numbers (DNs) given in the *WISE* level 1b images. The solid line is the best-fit Moffat function.

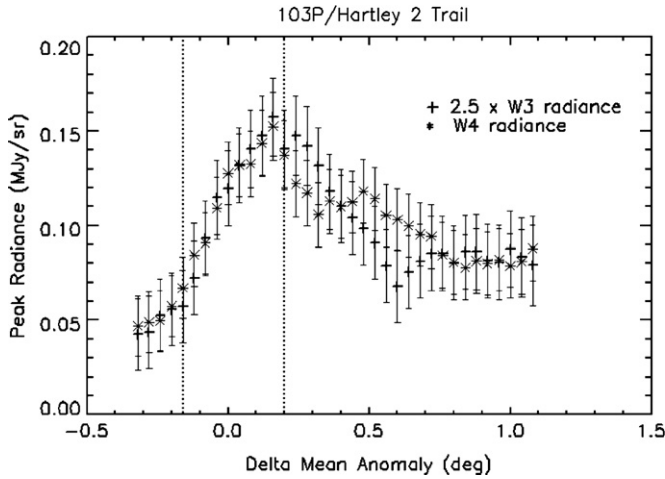
et al. 2004). The total mass loss in coma dust emission over the course of an orbit at the observed rate would be $\sim 2 \times 10^9$ kg, assuming 1 g cm^{-3} grain densities, or a rate of $\sim 7 \text{ kg s}^{-1}$, and is consistent with the values derived by Lisse et al. (2009).

Figure 3 shows the thermal and reflected light contributions from the dust to the measured flux, after the blackbody temperature fit. The nucleus signal, insignificant compared to the flux uncertainty, has also been removed from each band. A neutral-reflectance for the coma grains was used originally in

the computation for the combined thermal and reflected light dust contribution, owing to the near-solar $V-R$ coma colors found using the UH data set. However, 103P was reported to have a noticeably red coma at the time of the encounter (cf. Sitko et al. 2011), into the near-infrared, prompting us to show, for comparison, a coma reflected-light model with a reddening law based on Jewitt & Meech (1986) averaged out to $3.5 \mu\text{m}$, i.e., with dust reflectivity proportional to the wavelength to the approximately 0.2 power. The difference in the flux in W2

Table 4
Production Rates

Quantity	12 μm , 11"	12 μm , 22"	22 μm , 11"	22 μm , 22"	4.6 μm 11"/22"	R Band, 11"	R Band, 22"
m_g (kg)	8.1×10^{-13}		5.7×10^{-12}		1.8×10^{-16}
n_g	2×10^{16}	3×10^{16}	1.26×10^{16}	3.0×10^{16}	2.2×10^{20}
Σ_{dust} (m^{-2})	25	9	16	9	2.8×10^5
Q_{dust} (kg s^{-1})	0.6	0.5	2.7	3.2	8.8×10^4
$\langle N_{\text{CO}_2} \rangle$ (m^{-2})	1.75×10^{10}	1.32×10^{10}	...
$\langle N_{\text{CO}} \rangle$ (m^{-2})	1.70×10^9	1.26×10^9	...
Q_{CO_2} (s^{-1})	$3.5 (\pm 0.9) \times 10^{24}$
Q_{CO} (s^{-1})	$3.4 (\pm 2) \times 10^{23}$

**Figure 7.** W3 and W4 peak radiance (MJy sr^{-1}) vs. dMA (–dMA is ahead of the nucleus). The dotted lines bound the region of dMA where the trail radiance is contaminated by coma and tail emission. The maximum contamination is $\sim 10\%$ of the peak. Error bars are $\pm 1\sigma$.

between the neutral-reflectance and reddening law reflected-light scaling models is well within the reported photometric uncertainty. However, we chose to base further calculations on the original neutral-reflectance model based on our UH measurements.

4.1. 4 μm Anomalous Emission

Comet 103P's signal at 4.6 μm is about 0.4 mJy for an 11 arcsec aperture. The contribution from the thermal and reflected light of the coma was calculated to be 0.03 mJy (see Figure 3). A large excess flux in the 4.6 μm band suggests other contributions, possibly from emission lines. Recent work with the *Spitzer Space Telescope* IRAC instrument's similar band-pass centered at 4.8 μm shows the most likely candidate emission lines to be from CO and CO₂ lines (Pittichova et al. 2008; Reach et al. 2009). With a band pass of $\sim 1 \mu\text{m}$ (Wright et al. 2010), the 4.6 μm band contains both the CO (4.26 μm) and CO₂ (4.67 μm) mid-IR emission bands. Water emission lines appropriate for the measured temperatures ~ 200 K that reside in the mid-IR nearest the band pass cluster in the 2.7–33 and 5.5–6.3 μm range (Lisse et al. 2006; Woodward et al. 2007). The excess of the W2 flux is likely attributable to CO and CO₂ gas species, seen strongly in Hartley 2 by the *EPOXI* mission in 2010 (A'Hearn et al. 2011) and in previous apparitions (Weaver et al. 1994; Crovisier et al. 1997), while CO was found to be very underabundant in this comet and also in other comets as the dominant lines in the 4–5 μm spectral range by ISO (Crovisier et al. 1999; Woodward et al. 2008; Reach et al. 2009). Comet 103P's signal at 4.6 μm is about 0.4 mJy for an 11 arcsec

aperture. The contribution from the thermal and reflected light of the coma was calculated to be 0.03 mJy (see Figure 3).

Using the analysis techniques as presented in Pittichova et al. (2008, see their Equations (4) and (5)) we convert our 4.6 μm band fluxes to the number densities and production rates shown in Table 4, assuming a gas ejection velocity of 0.62 km s^{-1} . We scale the contribution from the two species for relative CO₂:CO abundance of roughly 10:1, similar to previous oppositions (Weaver et al. 1994; Bhardwaj & Raghuram 2011) to arrive at our relative production values. These represent the unique constraints on these species abundances for 103P at these distances. We have no available measurement of H₂O production near the time of the *WISE* observations, so we are unable to provide reliable constraints on the CO₂/H₂O ratio. However, extrapolation to 2010 May 10 from the production models as shown in Meech et al. (2011) and from rates observed on 2010 July 26 by Mumma et al. (2011) yields H₂O values $\sim 3 \times 10^{26}$ mol s^{-1} , suggesting that a derived CO₂/H₂O ratio would be on the order of a few percent, i.e., a factor of a few lower than seen at perihelion during the *EPOXI* encounter (Weaver et al. 2011) and for previous orbits (Weaver et al. 1994). The CO₂/H₂O ratio may possibly decrease as a function of heliocentric distance.

4.2. Dust Trail

The debris trail grain temperature (T_g) was calculated at each dMA from the ratio of the profile peak radiance in W3 to that in W4. The resulting temperatures were well represented over the range of trail distances ($2.26 \text{ AU} < r < 2.39 \text{ AU}$) by the relation $T_g = 303/r_h^{1/2} \pm 5 \text{ K}$, where r_h is given in AU and is in good agreement with the relation derived for *IRAS* trail comets with multiband detections (Sykes et al. 2005). As had been noted previously (Sykes et al. 1990), trail temperatures are significantly ($\sim 10\%$) higher than the local equilibrium blackbody temperature. Possible explanations have been discussed extensively (Sykes et al. 1990, 2005; Sykes & Walker 1992); the least complex explanation assumes that randomly orientated trail grains maintain a latitudinal temperature gradient across their surfaces due to their very low thermal inertia. The optical depth, τ , of the trail at the profile peak radiance was derived using the above relation for T_g and assuming that the trail grains emit blackbody radiation, and is plotted in Figure 9 for each dMA observed. A solid line is drawn through the means of the W3 and W4 results and represents the values of τ used to integrate the trail mass. The peak optical depths for 103P lie within the range of values $(0.3\text{--}16) \times 10^{-9}$ found for the 34 comets surveyed by *Spitzer* (Reach et al. 2007) and $(1.3\text{--}13.8) \times 10^{-9}$ found for the eight *IRAS* trails (Sykes & Walker 1992).

The width of the trail, as derived from the FWHM of the radiance profile and the *WISE* centric distance to the trail, is

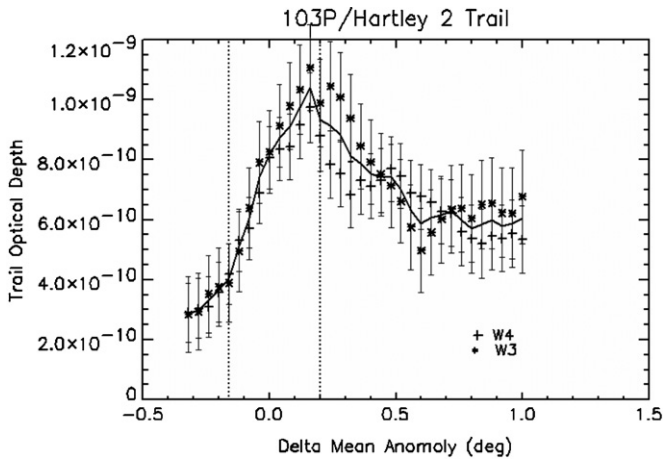


Figure 8. Optical depth at the peak of the radiance profile is shown. The solid line is drawn through the means of the W3 and W4 values, and defines the values used for the integration of trail mass. Error bars are $\pm 1\sigma$.

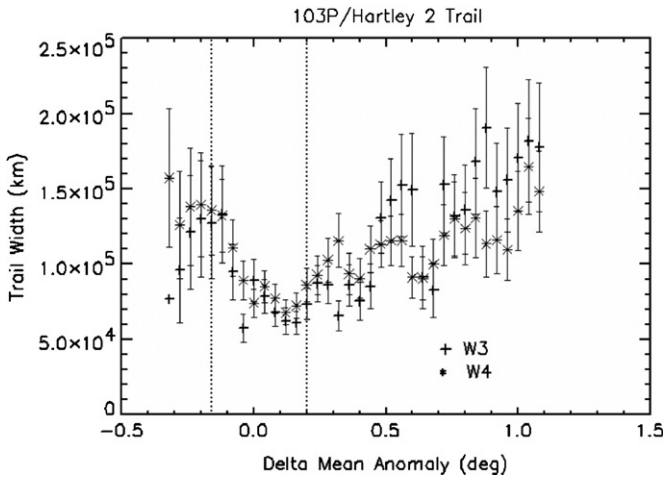


Figure 9. Width of the trail vs. dMA. The widths plotted are derived from the FWHM of the radiance profiles and are corrected for the distance from *WISE* to the trail. Error bars are $\pm 1\sigma$.

shown in Figure 9. The trail reaches a minimum width of $\sim 65,000$ km slightly behind the comet at $dMA = 0.12$, increases to a mean of $\sim 140,000$ km ahead of the nucleus ($-0.32 < dMA < -0.16$), and increases to $\sim 170,000$ km behind the nucleus at $dMA = 1.04$. Our value of $65,000$ km near the nucleus is higher than the $40,000$ km width found from $22\ \mu\text{m}$ *Spitzer* photometric imaging (Lisse et al. 2009), however, this probably is due to the “contamination/averaging” previously discussed in the Analysis section. As such, it should be considered an upper limit to the trail width near the nucleus.

The velocity of trail grains can be deduced from measurements of the physical trail width (Sykes et al. 1986). Applying their Equations (1) and (2) to the above measurements of the width of the 103P trail and assuming isotropic perihelion emission, we derive a mean grain velocity relative to the nucleus of $3.6 \pm 1.4\ \text{m s}^{-1}$. This is in good agreement with a radial velocity dispersion of the grains of $4\ \text{m s}^{-1}$ from preliminary radar observations (Harmon et al. 2010).

The grain size is the final element required to calculate the mass of the observed debris trail. Whereas we cannot measure the grain size directly, we can measure the ratio $\beta = \text{Force due to radiation} / \text{Force due to solar gravity} = 1.14 \times 10^{-4} Q_{\text{pr}} \rho d^{-1}$ (Burns et al. 1979) where Q_{pr} = efficiency for radiation pressure

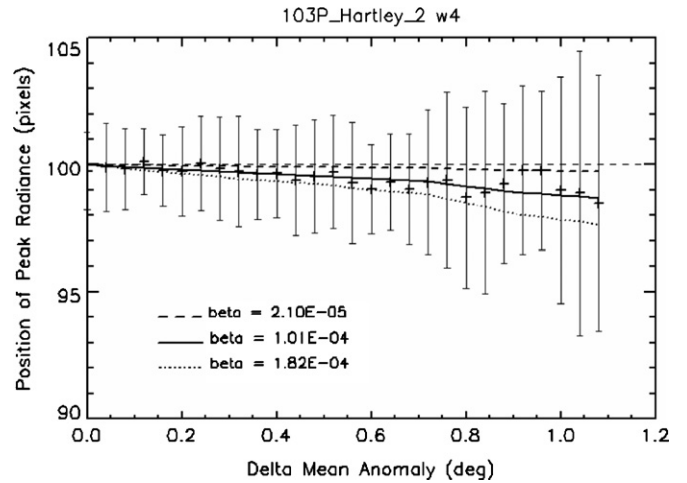


Figure 10. Positions of the peaks (+ symbols) of the radiance profile as a function of dMA. Error bars are $\pm 1\sigma$. The horizontal dashed line at pixel 100 is the position of the comet’s orbit, that is, the zero velocity and zero β syndyne. The solid line is the zero velocity syndyne for $\beta_{\text{min}} = 1.01 \times 10^{-4}$. The dotted and dark-dashed lines are the $\beta_{\text{min}} \pm 1\sigma$ values.

≈ 1.0 for grains with $d \gg \lambda$, ρ = grain density (g cm^{-3}) and d = grain diameter (cm). Trail grains ejected from the nucleus experience a reduced force of gravity due to solar radiation pressure that changes the semimajor axis of their orbits from that of the parent comet (Sykes & Walker 1992). Figure 10 is a plot of the position of the peaks of the trail radiance profiles behind the nucleus versus dMA. The horizontal dashed line at pixel 100 is the position of the comet’s orbit, that is, the zero velocity and zero β syndyne. Equation (2) from Sykes & Walker (1992) was used to calculate the positions of trail grains for a large set of β and find β_{min} giving a minimum χ^2 fit to the measured peak positions. The solid line is the zero velocity syndyne for $\beta_{\text{min}} = 1.0 \times 10^{-4}$. The dotted and dashed lines are for $\beta_{\text{min}} \pm 1\sigma$. In the case where $\rho = 1.0\ \text{g cm}^{-3}$, $\beta_{\text{min}} = 1.0 \times 10^{-4}$ corresponds to a grain diameter = $1.1\ \text{cm}$, a diameter typical of dust trail particles (Sykes & Walker 1992; Reach et al. 2007), comet-associated meteoroids, and the icy grains imaged around the nucleus of 103P by the *EPOXI* spacecraft (A’Hearn et al. 2011). The $\pm 1\sigma$ β_{min} values of 1.82×10^{-4} and 2.10×10^{-5} correspond to grain diameters of 0.63 and $5.4\ \text{cm}$, respectively. Note that the density for large particles may be $\sim 0.4\ \text{g cm}^{-3}$ (Richardson et al. 2007) and thus the diameter of the particles may be 2–3 times larger. This result agrees with preliminary radar Doppler spectra that show a broadband echo from $> 1\ \text{cm}$ diameter grains (Harmon et al. 2010).

Our solution for β also gives an estimate of the age of the trail particles. Grains observed at $dMA = 0.32$ were emitted approximately 6.5 years (1 period) ago, while the grains observed most distant from the nucleus were emitted about 21 years (3.3 periods) ago.

103P is observed to have a significant trail of debris leading the comet. The minimum particle diameter ahead of the comet is that for which the effect of ejection velocity balances that of radiation pressure (Sykes et al. 1986; Lisse et al. 2004). Setting the above referenced Equation (2) (Sykes & Walker 1992) to zero and solving for β yields an estimate of the maximum value of $\beta = 1.91 \times 10^{-4}$ with corresponding minimum diameter of $0.59\ \text{cm}$ in the leading portion of the trail (assuming $\rho = 1$).

The mass of the observed trail is estimated from the optical depth, peak normalized profile area, and β . If we assume that all the grains in a given field of view (FOV) are at the same

temperature, then the optical depth, τ , is just the fraction of the FOV area at the trail that is filled by the projected areas of the grains, thus $\tau = \pi n R^2 / (\alpha \Delta)^2$ where n is the number of grains of radius R within the FOV defined by the angular size of the pixel, α , and the *WISE* centric distance, Δ . The mass of the grains within the FOV is just the mass of a single spherical grain times the number of grains, that is, $M_{\text{FOV}} = 4/3 \pi R \rho \tau (\alpha \Delta)^2$ where ρ is the grain density. In terms of β and the *WISE* W4 pixel size of 5.5 arcsec this reduces to $M_{\text{FOV}} = 1.218 \times 10^{13} \tau \Delta^2 \rho / \beta$ (g pixel $^{-2}$) where Δ is in AU. The total mass of the observed trail is then $M_{\text{total}} = W M_{\text{FOV}}$, where the width of the trail, W , in pixels is the area under the peak-normalized radiance profile, and L is the length of the trail, in pixels, observed on the sky. In practice, since τ , Δ , and W vary with sky position, M_{total} was calculated as an integral along the trail length. Also, the portions of the trail ahead and behind the comet were treated separately, since each had a different value of β . This yields 3×10^{12} g for the total mass of the trail leading the nucleus, and 4×10^{13} g for the longer trail section in the following portion, assuming $\rho = 1 \text{ g cm}^{-3}$. As previously noted (Reach et al. 2007), these values are lower limits to the trail mass because the full extent of the trail may not have been measured and larger grains may be present. The mass of the 103P trail is within the range of 3×10^8 to 3×10^{11} kg found for the eight trails observed by *IRAS* (Sykes & Walker 1992) and the range 4×10^7 to 9×10^{10} kg found for six *Spitzer* comets “with reasonable total mass estimates” (Reach et al. 2007).

Given the age of the grains calculated above, we estimate the average mass loss rate of 103P meteoroids over 21 years is 62 kg s^{-1} into the trail, which is greater by a factor of nine than the average orbital mass loss rate of $\sim 7 \text{ kg s}^{-1}$ into coma dust. This rate would amount to the loss per orbit of $\sim 1.5\%$ of the mass of the comet’s nucleus, assuming a density near 1 g cm^{-3} . Our value falls within the range of $4\text{--}250 \text{ kg s}^{-1}$ found for the *IRAS* trails, but outside the range of $0.2\text{--}36 \text{ kg s}^{-1}$ for the *Spitzer* trails. If we exclude 2P/Encke and 29P/Schwassmann–Wachmann 1 from the *IRAS* list, then the *IRAS* range of mass loss rates for more typical JFC trails is $4\text{--}20 \text{ kg s}^{-1}$, totally within the range measured by *Spitzer*, and our value of 62 kg s^{-1} for 103P is now well outside that range.

5. CONCLUSIONS

Observed by *WISE* at a unique time and heliocentric distance relative to its spacecraft encounter, 103P exhibits interesting characteristics for a short period comet. These data help constrain behavior over a longer period than the immediate pre- and post-encounter observations, and the simultaneous and near-simultaneous data presented here in combination yield the following results.

1. The extracted thermal flux from the nucleus indicates a mean diameter consistent with encounter measurements for the geometric mean radius of the body. This is as expected, considering the ~ 38 hr span over which the 18 *WISE* images were taken covers two full rotations of the 18 hr period of the comet. The fitted η values of ~ 1.2 are marginally higher than seen in other data sets.
2. The CO_2 (and CO) detected in W2 excess is consistent with CO_2 production rates on the order of 3.5×10^{24} molecules per second on 2010 May 10 at $r_h = 2.3$ AU. This is lower than the estimated production seen near perihelion, at $r_h = 1.06$ AU (Meech et al. 2011; Weaver et al. 2011), but still plays a significant role in the cometary activity at these distances.

3. A blackbody curve fit to the observed thermal flux from coma dust particles yields dust temperatures near the temperature for a blackbody at the comet’s 2010 May 10 heliocentric distance, 2.3 AU, and is consistent with large-particle ($> 11 \mu\text{m}$) dust dominance.
4. The extended trail’s extracted peak radiance locations were fit with an array of dust β parameter values. The best-fit values are consistent with large-particle sizes on the order of centimeters, and with a long-lived dust trail that is extant for periods greater than a single orbit.
5. The total trail mass and mass production rate coupled with a small nuclear diameter indicate that 103P’s refractory component dominates its mass.

The *WISE* 103P data provide meaningful values for each observable that the data provided and serve as a good example of what the *WISE* cometary body data set is capable of yielding on individual bodies, and for statistical comparisons of the over 120 comets observed in the sample.

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